

# A Conjecture on Hermite Constants

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**Abstract.** As of today, the Hermite constants  $\gamma_n$  are only known for  $n \in \{1, 2, 3, 4, 5, 6, 7, 8, 24\}$ .

We noted that the known values of  $(4/\gamma_n)^n$  coincide with the values of the minimal determinants of any  $n$ -dimensional integral lattice when the length of the smallest lattice element  $\mu$  is fixed to 4.

Based on this observation, we conjecture that the values of  $\gamma_n^n$  for  $n = 9, \dots, 23$  are those given in Table 2.

We provide a supporting argument to back this conjecture. We also provide a provable lower bound on the Hermite constants for  $1 \leq n \leq 24$ .

## 1 The Observation

Hermite constants  $\gamma_n$  determine how short a lattice element can be.  $\gamma_n$  is defined as follows: Let  $\mathcal{L}$  be a lattice in Euclidean space  $\mathbb{R}^n$  with unit co-volume. Denoting by  $\lambda_1$  the minimal length of a nonzero element of  $\mathcal{L}$  we denote by  $\sqrt{\gamma_n}$  the maximal value of  $\lambda_1$  over all such lattices  $\mathcal{L}$ . Table 1 gives all known values of  $\gamma_n^n$ :

$n$	1	2	3	4	5	6	7	8	9...23	24
$\gamma_n^n$	1	$\frac{4}{3}$	2	4	8	$\frac{64}{3}$	64	$2^8$	?	$2^{48}$

**Table 1.** All known values of  $\gamma_n^n$

During the implementation of post-quantum encryption algorithms, we noted that the eight ratios between the first known Hermite constant values and  $2^n$  coincide with the values  $\alpha_{2,n}$  of the minimal determinants of any  $n$ -dimensional integral lattice when we fix the length of the smallest lattice element to  $\mu = 2$ . Because this coincidence does not hold for  $n = 24$  we further investigated what happens for higher  $\mu$  values.

It turns out that all 9 known values of  $(4/\gamma_n)^n$  coincide with  $\alpha_{4,n}$ , which is the equivalent of  $\alpha_{2,n}$  when  $\mu = 4$ . The quantities  $\alpha_{\mu,n}$  are known for  $\mu \in \{2, 3, 4\}, n \geq 0$  [CS99] and are lattice-related, which makes the following conjecture reasonably likely:

Conjecture 1. For  $1 \leq n \leq 24$ :

$$\gamma_n^n = \frac{4^n}{\alpha_{4,n}}$$

where  $\alpha_{4,n}$  is the minimal determinant of any  $n$ -dimensional 4-norm integral lattice.

If Conjecture 1 is true then the values of  $\gamma_n^n$  for  $n = 9, \dots, 23$  would be those given in **green** in Table 2.

In addition, our estimates coincide with two other values for  $n = 9, 10$  given (with an apparently incomplete proof [Cox47]) in [Cha46].

$n$	$\alpha_{2,n}$	$\alpha_{4,n}$	$2^n$	$4^n$	$2^n/\alpha_{2,n}$	$4^n/\alpha_{4,n}$	$\gamma_n^n$
1	2	4	$2^1$	$4^1$	$2^0$	$2^0$	$2^0$
2	3	12	$2^2$	$4^2$	$2^2/3$	$2^2/3$	$2^2/3$
3	4	32	$2^3$	$4^3$	$2^1$	$2^1$	$2^1$
4	4	64	$2^4$	$4^4$	$2^2$	$2^2$	$2^2$
5	4	128	$2^5$	$4^5$	$2^3$	$2^3$	$2^3$
6	3	192	$2^6$	$4^6$	$2^6/3$	$2^6/3$	$2^6/3$
7	2	256	$2^7$	$4^7$	$2^6$	$2^6$	$2^6$
8	1	256	$2^8$	$4^8$	$2^8$	$2^8$	$2^8$
9	2	512	$2^9$	$4^9$	$2^8$	$2^9$	$2^9$
10	3	768	$2^{10}$	$4^{10}$	$2^{10}/3$	$2^{12}/3$	$2^{12}/3$
11	2	972	$2^{11}$	$4^{11}$	$2^{10}$	$2^{20}/3^5$	?
12	1	729	$2^{12}$	$4^{12}$	$2^{12}$	$2^{24}/3^6$	?
13	2	972	$2^{13}$	$4^{13}$	$2^{12}$	$2^{24}/3^5$	?
14	1	768	$2^{14}$	$4^{14}$	$2^{14}$	$2^{20}/3$	?
15	1	512	$2^{15}$	$4^{15}$	$2^{15}$	$2^{21}$	?
16	1	256	$2^{16}$	$4^{16}$	$2^{16}$	$2^{24}$	?
17	1	256	$2^{17}$	$4^{17}$	$2^{17}$	$2^{26}$	?
18	1	192	$2^{18}$	$4^{18}$	$2^{18}$	$2^{30}/3$	?
19	1	128	$2^{19}$	$4^{19}$	$2^{19}$	$2^{31}$	?
20	1	64	$2^{20}$	$4^{20}$	$2^{20}$	$2^{34}$	?
21	1	32	$2^{21}$	$4^{21}$	$2^{21}$	$2^{37}$	?
22	1	12	$2^{22}$	$4^{22}$	$2^{22}$	$2^{42}/3$	?
23	1	4	$2^{23}$	$4^{23}$	$2^{23}$	$2^{44}$	?
24	1	1	$2^{24}$	$4^{24}$	$2^{24}$	$2^{48}$	$2^{48}$

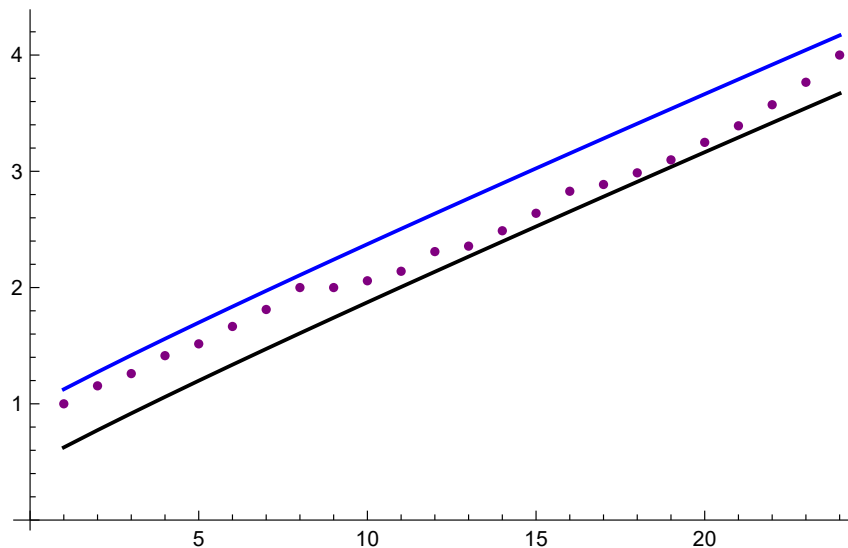
**Table 2.** Values of  $\gamma_n^n, \alpha_{2,n}, \alpha_{4,n}$ , the coinciding values are shown in **blue** and the mismatch for  $\mu = 2$  is shown in **red**. The two values of [Cha46] are shown in **magenta**. The conjectured values are given in **green**.

The observation fits with the known bounds of  $\gamma_n$  (the classical  $\Gamma$ -function one and [WCW19]). Namely:

$$\frac{n}{2\pi e} < \gamma_n = \frac{4}{\sqrt{\alpha_{4,n}}} \leq \frac{2}{\pi} \Gamma\left(2 + \frac{n}{2}\right)^{2/n} < \frac{n}{8.5} + 2 \text{ for } 1 \leq i \leq 24$$

The conjectured values that we get for  $\gamma_n^n$  are remarkably close to the line  $(n-1)/8+1$  but this is unfortunately incompatible with the (crude) upper bound having a slope of  $1/8.5$  as both lines meet at  $n = 153$ . The next “natural” line having a slope  $< 1/8.5$  and lower bounding all the values we found is  $(n-1)/9+1$ . Given that from 17 and on we see a takeoff from  $(n-1)/9+1$ , if this takeoff is indeed preserved for higher values we may consider  $(n-1)/9+1$  as a plausible candidate lower-bound. We also note that for all 24 known and conjectured values:

$$0 < \Gamma(2 + \frac{n}{2})^{2/n} - \gamma_n < \frac{1}{2}$$

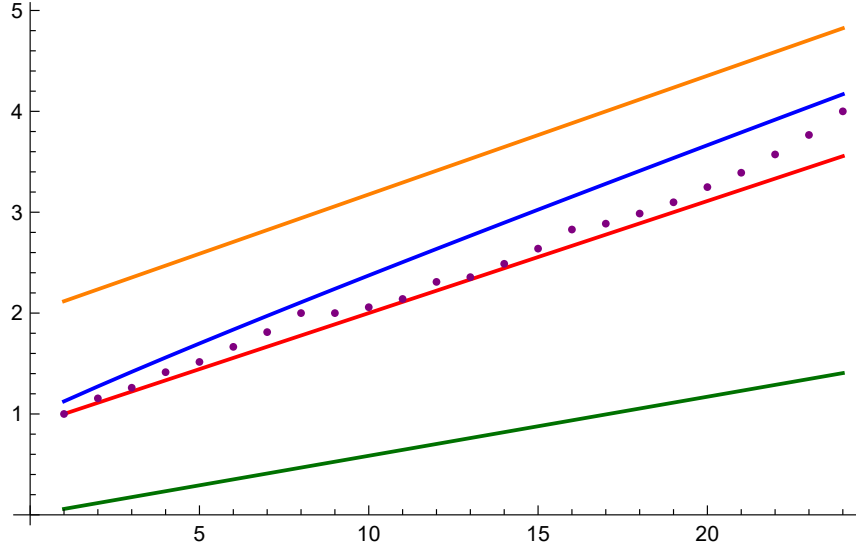


**Fig. 1.** The conjectured values are shown in violet. The blue line is the proved  $\Gamma$ -function bound. The black line is the  $\Gamma$ -function bound minus  $1/2$

## 2 Motivation

The outline of our reasoning is the following: we depart from the observation that the problem of computing Hermite constants and the problem of computing smallest determinants are equivalent. However, to our knowledge solutions to the smallest determinant problem are only known for integral lattices and  $\mu = 2, 3, 4$ .

We conjecture that the two problems are still equivalent even when one restricts the minimal determinant problem to even integral lattices. First we restate some well known definitions.



**Fig. 2.** The conjectured values are shown in violet. The red line is  $(n - 1)/9 + 1$ . The orange line is the proved upper bound  $n/8.5 + 2$ , the blue line is the proved  $\Gamma$ -function bound. The green line is the proved lower bound  $n/(2\pi e)$ .

**Definition 1 (Hermite function).** Let  $\mathcal{L}_n$  be the set of all lattices with dimension  $n$  and  $L \in \mathcal{L}_n$ . The Hermite function is defined as:

$$\gamma(L) = \frac{\mu}{\det(L)^{\frac{1}{n}}}$$

where  $\mu = \min(L)$  is the length of the smallest nonzero element of  $L$ .

A simple property of the Hermite function is the invariance under scaling.

**Theorem 1 (Invariance under scaling).** Let  $c \in \mathbb{R}_{>0}$  and  $\gamma(L)$  be the Hermite function. It holds that:

$$\gamma(cL) = \gamma(L)$$

*Proof.* This follows directly from the definition of the Hermite function. □

**Definition 2 (Hermite constant).** Let  $\mathcal{L}_n$  be the set of all lattices with dimension  $n$  and  $L \in \mathcal{L}_n$ . The Hermite constant  $\gamma_n$  is defined as:

$$\gamma_n = \sup\{\gamma(L) | L \in \mathcal{L}_n\}$$

We now have all the necessary definitions and theorems but let us first give the intuition behind the equivalence theorem. To find a Hermite constant  $\gamma_n$ ,

imagine that one starts with an arbitrary  $L \in \mathcal{L}_n$  with:

$$\gamma(L) = \frac{\mu}{\det(L)^{\frac{1}{n}}}$$

The goal is now to modify  $L$  in a way that increases  $\gamma(L)$  up to  $\gamma_n$ . Clearly, this can only be achieved by increasing  $\mu$  or by decreasing  $\det(L)$ .

From the invariance under scaling we know that it is allowed to either fix  $\det(L)$  and increase  $\mu$  or fix  $\mu$  and decrease  $\det(L)$ , one can always scale up or down in the end. Note that both approaches solve the same problem: maximizing the ratio. Solutions to both problems are thus equivalent. More formally:

**Definition 3** ( $L_1, L_2, L_3$ ). *Given the set  $\mathcal{L}_n$  of all lattices with dimension  $n \in \mathbb{N}$ . Let  $L_1 \in \mathcal{L}_n$  be:*

$$L_1 = \operatorname{argmax}_L \frac{\min(L)}{\det(L)^{\frac{1}{n}}}$$

Let  $L_2 \in \mathcal{L}_n$  be:

$$L_2 = \operatorname{argmax}_L \frac{1}{\det(L)^{\frac{1}{n}}}$$

Let  $L_3 \in \mathcal{L}_n$  be:

$$L_3 = \operatorname{argmax}_L \min(L)$$

where again  $\min(L) = \mu$  is the length of the smallest nonzero element of  $L$ .

**Theorem 2 (Equivalence).** *Let  $\mathcal{L}_n$  be the set of all lattices of dimension  $n$ ,  $\gamma(L)$  be the Hermite function,  $\gamma_n$  the Hermite constant of dimension  $n$  and  $L_1, L_2$  and  $L_3$  be defined as above. It holds that:*

$$\gamma_n = \gamma(L_1) = \gamma(L_2) = \gamma(L_3)$$

*Proof.* This follows directly from Theorem 1. □

## 2.1 The problem of integral lattices

The question remains why the sequences calculated from  $\alpha_{2,n}$  and  $\alpha_{4,n}$  differ for  $n \geq 9$ . This follows from the fact that the values of  $\alpha_{\mu,n}$  were computed with the restriction to integral lattices. This means that the determinant cannot get smaller than 1. Now a problem arises when  $\mu$  is fixed and  $\det(L) = 1$ . Under such circumstances we can no longer decrease  $\det(L)$  and  $\mu$  is fixed so the only scaling comes from the dimension (the root in the denominator).

This explains perfectly why  $\frac{2^n}{\alpha_{2,n}}$  works until  $n = 8$  but not after. Because for  $n = 8$  we hit  $\det(L) = 1$ . This also suggests that the results we obtained for  $\frac{4^n}{\alpha_{4,n}}$  are only optimal until  $n = 24$ , but also gives good reason to conjecture that until then they are in fact optimal. More importantly, this gives an indication how to find the next Hermite constants for  $n > 24$ : solve the problem of the minimal

determinant of  $n$  dimensional lattices with a bigger  $\mu$ . Whether  $\mu$  needs to be a power of two or just even is unclear. To be safe we choose even. This in turn allows us to formulate the more general conjecture:

*Conjecture 2.* For any  $n \geq 0$ :

$$\gamma_n^n = \frac{\mu^n}{\alpha_{\mu,n}}$$

where  $\alpha_{\mu,n}$  is the minimal determinant of any  $n$ -dimensional  $\mu$ -norm integral lattice and  $\mu$  is even and chosen big enough.

## 2.2 Lower bound on $\gamma_n$

If in fact it should turn out that neither Conjecture 1 nor 2 are true we note that the values from our observation still provide a provable lower bound on the Hermite constants for  $1 \leq n \leq 24$ .

**An open question:** It is clear that the restriction to integral lattices in the calculation of the  $\alpha_{\mu,n}$  will lead to problems when the determinant reaches its minimum. It can also be seen that the approach doesn't work for  $\mu = 3$ . We think that this also follows from this restriction to integral lattices. The question remains why this restriction does not lead to other problems and non-optimal solutions for the even/power-of-two lattices before the determinant limit is reached. In our reasoning the restriction to integral lattices is not needed but to our knowledge only solutions to the restricted problem are known.

**Reverse-engineering the constants?** Note that all known and conjectured Hermite constants up to  $\gamma_{24}^{24}$  can be written as a power of 2 that is sometimes divided by a power of three. We assume the conjecture to hold and denote this as

$$\gamma_n^n = \frac{4^n}{\alpha_{4,n}} = \frac{2^{v_n}}{3^{u_n}} = 2^{v_n - \log_2(3)u_n}$$

These powers feature a symmetry in value and position if we start from  $n = 0$ , see Table 3. This can be written as:

$$\gamma_{24-n}^{24-n} = 2^{48-4n} \gamma_n^n \text{ for } 1 \leq n \leq 24$$

Given that there is a symmetry around element  $n = 12$ , it might be possible that this symmetry is further preserved after  $n = 24$  in a way unknown so far, thereby revealing more information on higher value Hermite constants.

Let  $\delta = 2 - \log_2(3)$  and  $\ell(n) = n \log_2 \gamma_n = u_n - \log_2(3)v_n$ , and denote by  $\bar{n} = \lfloor 6 - \frac{n}{2} \rfloor$ . Then:

$$\lfloor \ell(n) \rfloor = \left\lfloor \frac{(n-4)^2}{8} \right\rfloor - |\bar{n} - 2| - \bar{n} + 8$$

This can be further refined because  $\ell(n) - \lfloor \ell(n) \rfloor$  equals  $-\delta$  when  $n = 2 \bmod 4$ . We can hence add those periodical values to further refine the estimator to:

$$\Delta_n = \left\lfloor \frac{(n-4)^2}{8} \right\rfloor - |\bar{n} - 2| - \bar{n} + 8 + \delta \left\lfloor \frac{n-3 \bmod 4}{3} \right\rfloor$$

$\Delta_n$  differs from  $\ell(n)$  only by  $5\delta - 2 \simeq 0.07518$  for  $n = 11, 13$  and by  $6\delta - 2 \simeq 0.49022$  for  $n = 12$ . As values range between 0 and 48, the relative order of magnitude of this discrepancy is  $\simeq 1\%$  of the observed phenomenon's scale for  $n = 12$  and even less for  $n = 11, 13$ .

It is possible to silence the discrepancies for  $n = 11, 13$  by adding to  $\Delta_n$  a further corrective term of

$$(5\delta - 2)(1 - \bar{n} + |1 - \bar{n}|)$$

and then tackle  $n = 12$  in a similar manner but this renders the formula more artificial and seems to hinder rather than clarify its nature. With such a correction the two quantities differ only by  $4\delta - 2 = 0.33985$  at one single point ( $n = 12$ ). Note that this formula might summarize the behavior for  $n \leq 24$  but not beyond  $n = 26$  because for  $n \geq 26$  the conjectured value of  $\lfloor \ell(n) \rfloor$  exceeds the  $\Gamma$  upper bound.

The relationship  $u_n = \lfloor \frac{5}{8}(v_n - \lfloor \ell_n \rfloor) \rfloor^1$  and the near-coincidence between  $u_n$  and  $(v_n - \lfloor \ell_n \rfloor)/2$  (that differ by one unit at  $n = 11, 12, 13$ ) clearly point to a structure yet to be explained. The following additional observations related to similar mathematical objects may provide further hints.

**Additional observations:** It is worthy noting that  $\lfloor n^2/8 \rfloor$  behavior is also deeply hidden in Table I page 594 of [CS82]. We first noticed that  $2^n \lambda_{24+n} = \lambda_n$  for  $n = 0, \dots, 23$ . Then defining

$$\frac{4^{n-24}}{\lambda_n} = \frac{2^{u'_n}}{3^{v'_n}} \text{ for } n = 25, \dots, 47$$

we noted that  $v'_n = 1$  when  $n \bmod 4 = 2$  and that  $f(n) = u'_n + 2v'_n - \lfloor \frac{(n-24)^2}{8} \rfloor$  presents three perfect linear regularities shown in Figure 4.

This provides a direct way to compute  $\lambda_n$ . A direct beautiful formula for  $\lambda_n$ , valid for  $n = 0, \dots, 47$ , is:

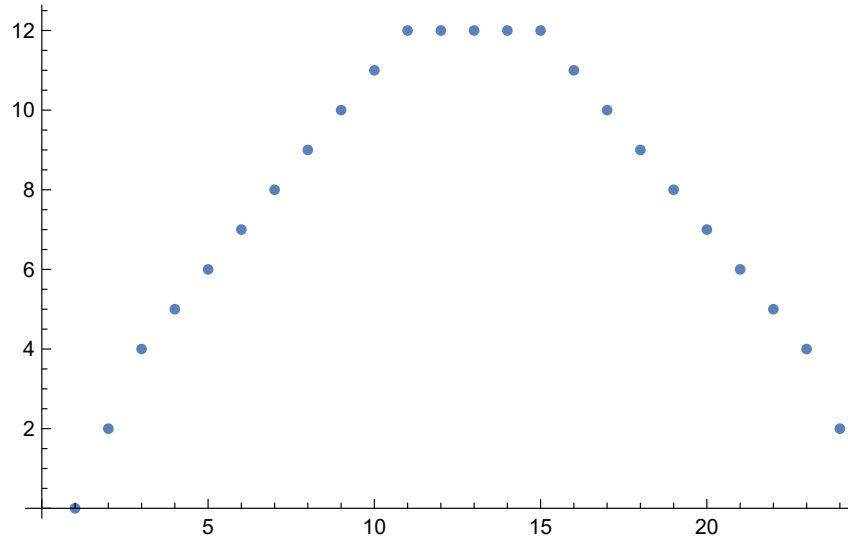
$$\lambda_n = \left(\frac{3}{4}\right)^{\lfloor \frac{n+1 \bmod 4}{3} \rfloor} \times 2^{|\bar{n}| - 2 + |\bar{n}| - \lfloor \frac{(\bar{n}-2)^2}{2} \rfloor - n \lfloor \frac{n}{24} - 1 \rfloor - 2} \text{ where } \bar{n} = 6 - \frac{n \bmod 24}{2}$$

A similar plateau appears in lattice  $K_n$  (Table II, page 601 of [CS82]) as shown in Figure 3 where  $u''_n$  and  $v''_n$  denote the powers of 2 and 3 found in  $\kappa_n$ .

<sup>1</sup> This simply stems from the fact that  $\lfloor \frac{5}{8} \lceil n \log_2(3) \rceil \rfloor = n$  for  $1 \leq i \leq 16$ .

$n$	$u_n$	$2n - v_n$
0		
1		2
2	1	2
3		5
4		6
5		7
6	1	6
7		8
8		8
9		9
10	1	8
11	5	2
12	6	
13	5	2
14	1	8
15		9
16		8
17		8
18	1	6
19		7
20		6
21		5
22	1	2
23		2
24		

**Table 3.** Symmetry in the powers of 2 and 3 in  $4^n/\alpha_{4,n}$ .



**Fig. 3.**  $u_n'' + 2v_n''$  for  $n = 1, \dots, 24$ .



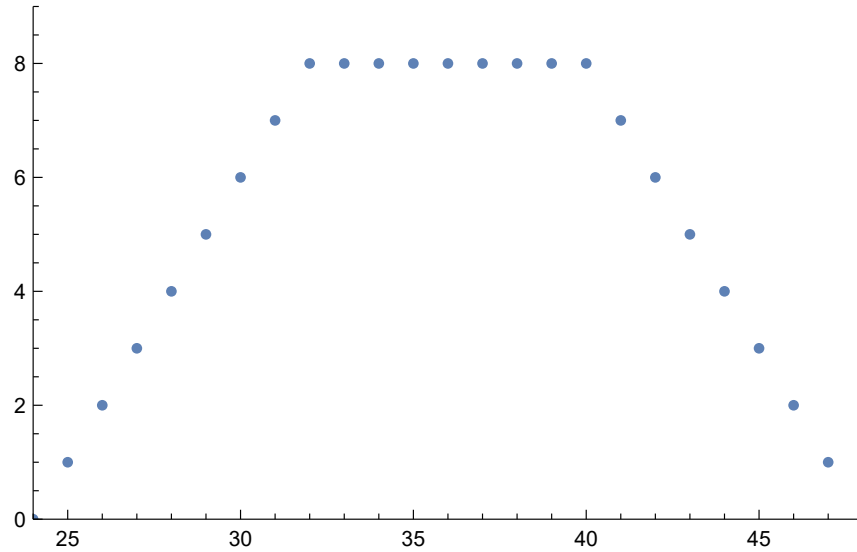


Fig. 4.  $u'_n + 2v'_n - \lfloor \frac{(n-24)^2}{8} \rfloor$  for  $n = 25, \dots, 47$ .

### 3 Conclusion

This note formulates conjectures relating to the values of Hermite constants for  $9 \leq n \leq 23$ . We explicit coherent regularities in the values of the conjectured constants and show that other lattice-related constants share similar features, notably a behavior governed by  $\lfloor \frac{(n-c_1)^2}{c_2} \rfloor$  for constants  $c_1, c_2$ .

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